

Improving Stability of Electronically Controlled Pressure Reducing Valves Through Gain Compensation

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ABSTRACT

This paper explains the root cause of instabilities which tend to arise in pressure reducing valves (PRVs) under low flow conditions. It was found that the loss of stability in PRVs is a direct result of an increase in the static valve/network gain as the valve position gets smaller, thus making pressure changes more sensitive to valve position adjustments. If the valve controller is tuned at medium valve openings characteristic of normal operating conditions, the increased gain at low valve openings can cause the control system to be too aggressive in its valve position adjustments leading to oscillations. The manuscript provides a mathematical derivation of the gain equation for a simplified pipe-PRV-pipe model. The obtained gain equation curve is then used to derive the formula for a gain compensator whose purpose is to keep the static gain constant across an entire range of permitted valve openings. A simplified network transient model is then used to recreate a real-life PRV instability event and show the remedial effects of the gain compensator.

Keywords: pressure reducing valve, instability, static gain, gain compensator.

INTRODUCTION

Energy and resource efficient operation of water distribution networks (WDNs) relies on an adequate choice of pressures at critical nodes within the network and an appropriate control of

24 these pressures such that they are maintained close to setpoints regardless of the changes in nodal
25 demands. It is up to engineering judgment to select the desired pressure values at controlled nodes
26 in the WDN bearing in mind that a minimum pressure, e.g. 15, is required at the tap to provide
27 an adequately high flow rate and prevent low pressure failures (Ghorbanian et al. 2015) whilst an
28 increase in pressure inside the pipes beyond the required minimum elevates the risk of pipe bursts
29 and leads to higher losses of water from the existing leaks (Thornton and Lambert 2005). An
30 appropriately designed pressure control system within a WDN will provide good setpoint tracking,
31 i.e. pressure values reflect the changes in the setpoint, good disturbance rejection, i.e. nodal
32 pressures are kept close to the setpoint regardless of demand changes, and additionally will not
33 cause large pressure transients in the system (stability). Although all three above requirements are
34 important, stability is crucial for the infrastructure since pressure transients put additional strain on
35 WDN components and, if high enough, can lead to pipe bursts.

36 Pressure in WDNs is often controlled by pressure reducing valves (PRVs). A PRV is a valve
37 combined with a controller in a form of a feedback control loop, as depicted in Fig. 1. The purpose
38 of this feedback controller is to maintain the outlet pressure at a constant value or to follow a
39 predefined trajectory irrespectively of inlet pressure and flow (demand) variations. The controller
40 receives an error signal e calculated as a difference between the downstream pressure setpoint H_d^{set}
41 and the measured downstream pressure H_d , and acts on this signal such that the error is minimized,
42 i.e. $H_d \approx H_d^{set}$. If $H_d < H_d^{set}$ the controller actuates the valve to increase the valve opening and
43 thus the outlet pressure. In the opposite case valve opening is decreased. The valve/WDN system,
44 also referred to as a plant, is subject to unknown disturbances in the form of demand changes,
45 which can be described as variations in orifice areas $A_{orifice}$ at the ends of demand nodes. If we
46 remove the feedback the system under study (i.e. the plant plus the controller) is called open-loop,
47 as opposed to full system with feedback which is referred to as closed-loop. This is a typical control
48 system and its properties, in particular stability, should be analyzed in the context of control theory
49 (e.g. Ogata 2010). From the point of view of the technology in which the controller is implemented
50 PRVs can be divided into two categories: hydraulically controlled valves and electrically controlled

51 valves. A hydraulic PRV is typically based on a pilot valve loop where the current outlet pressure
52 is compared with the set-point defined by the preload of a spring in the pilot valve (Prescott and
53 Ulanicki 2003). In addition to a hydraulic PRV two types of electronically controlled PRVs are
54 available: (a) a hydraulic PRV such as the one mentioned above in which the set-point is adjusted
55 by an electronic actuator and (b) a standard valve coupled with an electronic feedback controller
56 which directly actuates the valve position. Both valves in the engineering jargon are referred to as
57 electronically controlled PRVs.

58 Since PRVs are a crucial component of every WDN, pressure control with PRVs has been a
59 widely researched topic. Unfortunately, most of the existing publications on pressure control are
60 focused on steady-state behavior with omission of controller and plant dynamics. A good review of
61 the state of the art in this area is given by Vicente et al. (2015). Traditionally PRVs have been used to
62 maintain a desired pressure immediately downstream of the valve but recently, the so called remote
63 real time control (RRTC) of water distribution networks (WDNs) where the controlled pressure is at
64 a remote node within a network has been of interest to a number of research groups. Campisano et al.
65 (2012) proposed a simple linear controller for real time control (RTC) of a motorized valve using
66 a steady-state model of a WDN. Similar idea was used for planning of remote real time controlled
67 PRVs for the Oppegård municipality in Norway (Berardi et al. 2015). In this case steady-state
68 WDN model incorporated pressure dependent background leakage. The same authors developed
69 a software module called WNetXL to support their methodology (Giustolisi et al. 2015). The
70 methodology was further refined by Campisano et al. (2016) through inclusion of the criteria for
71 target node selection, allowing multiple valves to control pressure at the same target node and by
72 considering controller calibration methods as well as wireless communication protocols. Page et al.
73 (2016) proposed an innovation to the design of real-time proportional controllers assuming that
74 the information about the flow through the PRV is available. Recently there has been a positive
75 trend to consider and understand dynamic behavior of valve/WDN systems. As an example, Creaco
76 et al. (2017) performed a simulation study of a RTC valve model combined with an unsteady flow
77 simulator with pulsed nodal demands. In this particular simulation study the results showed some

78 pressure and flow oscillations but the authors concluded that the steady-state model was a good
79 approximation for the design of the RTC schemes and PRV stability was not discussed in the paper.

80 Dynamic behavior of PRVs was also tested experimentally, although none of the empirical
81 research based publications known to authors deal with PRVs during unstable operation, only with
82 stable dynamic response to disturbances. Before we move on to further discussions on PRVs'
83 dynamic behavior it is worth pointing out that PRVs alike other dynamic systems may show an
84 oscillatory response (e.g. to incoming pressure waves) but this is not to be confused with unstable
85 response. Stable response may be sinusoidal as long as its amplitude decays in time whilst instability
86 manifests itself with a growing output for a bound input, e.g. a sinusoid of an amplitude growing
87 in time. As pointed out in Janus and Ulanicki (2017) PRV instability is likely to cause pressure
88 transients but itself is a sole result of dynamic properties of the system, as shall be discussed in
89 the later sections of this manuscript. Instabilities are not caused by pressure transients and in fact
90 can be shown (numerically) to occur in rigid-column simulations where pressure transients are
91 not modeled. Dynamic behavior of PRVs was first analyzed by Brunone and Morelli (1999) who
92 measured the response of an automatic control valve to changes in the position of the valve at
93 the confluence node downstream of the PRV. The results show higher pressure oscillations under
94 lower flows than higher flows which supports the findings of this manuscript. The publication also
95 introduces a technique to obtain the flow-rate curve of a valve through unsteady-state tests, which
96 offers better estimation accuracy, especially under lower flows. Meniconi et al. (2015) measured
97 pressure waves produced during partial a partial closure and opening of a valve downstream of a
98 PRV and initially analyzed the waves using a wavelet transform. Experimental studies on dynamic
99 response of PRVs were taken forward in Meniconi et al. (2016) and Meniconi et al. (2017) who
100 performed further laboratory experiments to monitor the behavior of their PRV under different flow
101 conditions. The authors observed pressure oscillations resulting from the movement of the valve
102 element but since the tested range of flows was quite narrow and, as explained and demonstrated
103 through simulation in Ulanicki and Skworcow (2014) and later in Janus and Ulanicki (2017), PRVs
104 tend to become unstable under sufficiently low flows, instability was not observed.

105 Since no article has yet discussed specifically the matter of PRV stability whilst the engineering
106 community needs to deal with often disastrous, as shall be briefly shown below, incidents of PRV
107 instability in WDNs, there is an urgent need to understand why such incidents occur and how
108 they can be prevented. Understanding stability is especially important where the valve element
109 is directly actuated by an electronic controller in RTC schemes. There are two potential reasons
110 for the instability in control schemes in general: (a) increase in the value of the system gain (here
111 increase of the static valve/WDN gain due to the nonlinearity of the valve/WDN system against
112 flow) and (b) increase in the value of the delay in the system due to long actuation times or too long
113 a sampling period. Whilst the latter can be avoided by proper system design, this paper focuses
114 on the former, which is caused solely by the property of the valve and the network. The aim of
115 this paper is to explain this phenomenon in a theoretical manner and derive a precise analytical
116 relationship between the valve opening and the static gain of the valve/WDN system.

117 How severe a PRV instability can be is demonstrated in one such instability event which occurred
118 in a large-scale pressure control scheme installed in a WDN of one of the major cities in the UK. The
119 above-mentioned instability event is illustrated in Fig. 2 which shows how the valve outlet pressure
120 and valve position began to oscillate under low valve openings (and hence low flow conditions)
121 whilst the PRV managed to regain stability later in the day when the flow increased and the valve
122 settled into a larger opening. The above event incurred serious financial and environmental costs
123 as it caused multiple pipe burst across the network resulting in the loss of water and the loss of
124 service to the consumers. Although, in this case the PRV under study was controlled electronically
125 by a programmable logic controller (PLC) the root-cause of instability is common for all types of
126 PRVs, only the remedy will be different for electronically controlled and mechanically controlled
127 PRVs and shall be subjects of future research papers.

128 The paper is organized as follows. First, the effects of the static valve/WDN gain on feedback
129 control system stability are demonstrated using a classical control root-locus technique applied to
130 a lumped ordinary differential equation (ODE) model designed to qualitatively represent transient
131 behavior of a WDN. Second, a theoretical formula for the static gain of a simplified valve/WDN

132 system consisting of an upstream pipe, PRV, and a downstream pipe is derived using implicit
 133 function theorem applied to static network equations. Next, 15-min measurements of flows and
 134 pressures obtained from a real network are used to calibrate the simplified hydraulic model such
 135 that both the model and the real network exhibit the same static hydraulic properties. The calibrated
 136 model is then used to generate the static valve/WDN gain vs. valve position curves for different
 137 values of the pressure dependency coefficient α in the pressure-dependent demand model. The
 138 curve for a chosen value of α is then used to derive the model of the static gain compensator, the
 139 purpose of which is to keep the static gain at a constant value for all permitted valve openings.
 140 Finally, the simplified network transient model is used to recreate a real-life PRV instability event
 141 briefly described above and to show the remedial effects of the gain compensator.

142 **EFFECT OF STATIC GAIN ON STABILITY OF A FEEDBACK SYSTEM**

143 The aim of this section is to explain the effect of the static gain of the open-loop valve/WDN
 144 system on the stability of the closed-loop system bearing in mind that a water distribution network
 145 equipped with a PRV is an example of a closed-loop control system (e.g. Ogata 2010) and can be
 146 represented in a simplified block-diagram form as in Fig. 1. The physical meaning of the static
 147 gain of the valve/WDN system in the context of pressure control is the ratio between the change
 148 in the downstream (outlet) pressure H_d and the change in the valve position x in steady-state, i.e.
 149 $K(x) = dH_d/dx$. The static gain depends on the valve opening (operating point) x signifying that
 150 the system is nonlinear. In broader terms, the valve/WDN system is nonlinear and distributed, i.e.
 151 it is described by nonlinear partial differential equations (PDEs). The approach usually adopted
 152 by control engineers in such a situation is to design a controller for a given operating point x
 153 using an approximate lumped model described by a linear time invariant (LTI) ordinary differential
 154 equation (ODE) or a system of LTI ODEs. Subsequently, the controller parameters are adapted to
 155 changes in the operating point, which is referred to as gain scheduling. A generic LTI ordinary
 156 differential equation has the the following form:

$$157 \quad a_n y^{(n)} + a_{n-1} y^{(n-1)} + \dots + a_1 y^{(1)} + a_0 y = b_m u^{(m)} + b_{m-1} u^{(m-1)} + \dots + b_1 u^{(1)} + b_0 u \quad (1)$$

158 The above equation can be then transformed from differential form in time domain into an algebraic
 159 form in the new Laplace s domain using the Laplace transform (e.g. Hazewinkel, M. 1994) and
 160 rearranged to produce the ratio between the output and the input. This ratio is called the transfer
 161 function and is shown in Eq. 2.

$$162 \quad P(s) = \frac{Y(s)}{U(s)} = \frac{b_m s^m + b_{m-1} s^{m-1} + \dots + b_1 s + b_0}{a_n s^n + a_{n-1} s^{n-1} + \dots + a_1 s + a_0} \quad (2)$$

163 The polynomials in the numerator and the denominator can be then factorized using general algebra

$$164 \quad P(s) = \frac{b_m (s - z_1) (s - z_2) \dots (s - z_{m-1}) (s - z_m)}{a_n (s - p_1) (s - p_2) \dots (s - p_{n-1}) (s - p_n)} \quad (3)$$

165 where the roots of the numerator, $z_1 \dots z_m$, are called zeros of the transfer function and the roots of
 166 the denominator, $p_1 \dots p_n$, are referred to as the poles of the transfer function. The zeros and the poles
 167 are either real or appear in complex conjugate pairs. The poles and zeros completely characterize
 168 Eq. 1 and hence provide a complete description of the system dynamics. In particular the poles
 169 directly define the homogeneous solution of the differential equation, $y_h(t) = \sum_{i=1}^{i=n} C_i e^{p_i t}$, where C_i
 170 denotes a constant coefficient next to the i -th exponential term of the solution. In other words the
 171 poles of the transfer function are the eigenvalues of the differential equation. Clearly, if the poles
 172 have negative real parts the homogeneous solution decays to zero and the system is stable. If at
 173 least one pole has a positive real part the homogeneous solution diverges to infinity and the system
 174 is unstable. The poles and the zeros can be plotted in the complex plane where the horizontal axis
 175 represents the real parts $\Re(s)$ and the vertical axis represents the imaginary parts $\Im(s)$. The transfer
 176 function of the closed-loop system displayed in Fig. 1 is $G_{closed}(s) = \frac{P(s) C(s)}{1 + P(s) C(s)}$, where $P(s)$
 177 is the plant transfer function and $C(s)$ is the controller transfer function. The expression in the
 178 numerator is called open loop transfer function $G_{open}(s) = P(s) C(s)$ and represents the transfer
 179 function of the system in the absence of feedback. If the parameters of the plant or the controller
 180 change so do the poles of the closed-loop system. If only one parameter is being varied, typically
 181 the static gain of the open loop system, the poles of the closed-loop system will move along curves

182 in the complex $\Re(s) - \Im(s)$ plane and such a graphical method of stability analysis is called
 183 root-locus (Evans 1950). The steady-state gain of the transfer function in Eq. 2 equals b_0/a_0 and
 184 corresponds to the static gain of the valve/WDN system $K(x) = dH_d/dx$ at a given operation point.
 185 As the gain increases these, hopefully originally stable, poles travel along the root-loci and, at
 186 some point, depending on the properties of the plant and the controller, may cross the imaginary
 187 axis where the closed-loop system becomes unstable. For the purpose of demonstrating how the
 188 static open loop gain affects the stability of the closed-loop system using classical control theory
 189 the authors developed examples of conceptual lumped models describing the downstream pressure
 190 head H_d in response to the changes in the valve position x . The first model is depicted in Fig. 3
 191 and was formulated based on the physical understanding of one-dimensional unsteady flow. It is
 192 hypothesized that the response of downstream pressure to the changes in valve position can be
 193 described as a superposition of two responses: (a) inertial response describing acceleration and
 194 deceleration of the mass of water under an assumption of zero compressibility, i.e. rigid column
 195 model and (b) the transient response, i.e. pressure oscillations due to water compressibility effects.
 196 It is also hypothesized that whilst the inertial response is a function of valve position, the transient
 197 response is a function of the change of valve position, i.e. the rate of valve closure/opening.
 198 Provided that the valve closes with a sufficiently small rate, i.e. $dx/dt \approx 0$ the transient response
 199 will be close to zero and the output will exhibit a predominantly inertial character, which is
 200 conceptually in agreement with what is observed in practice. The inertial response is modeled
 201 with a first order lag element $\frac{K_1}{\tau s + 1}$ whilst the oscillatory response is described with a canonical
 202 second order system $\frac{K_2 \omega_n^2}{s^2 + 2 \zeta \omega_n s + \omega_n^2}$, where τ denotes the system time constant, ω_n denotes the
 203 natural frequency and ζ is the damping ratio (see e.g. Ogata 2010). As an alternative conceptual
 204 lumped model of a valve/WDN system we can use a Maxwell model of a mass supported in parallel
 205 by a spring and spring with a damper (see Fig. 4). The Maxwell model exhibits similar behavior to
 206 the inertial-oscillator model described above - oscillations superimposed on slow inertial response.
 207 The transfer function of the Maxwell model between the force applied to mass m and the mass

position is as follows:

$$P_{Maxwell}(s) = \frac{s + \frac{E_2}{\eta}}{s^3 + \frac{E_2}{c}s^2 + \frac{E_1+E_2}{m}s + \frac{E_1 E_2}{m \eta}} \quad (4)$$

where E_1 is the spring constant of the left spring, E_2 is the spring constant of the right spring, m is the mass of the body supported by the springs and η is the viscous friction coefficient of the damper. To visualize the dynamic behavior of the Maxwell model to a unit step change in the external force the following values of model parameters were assumed: $m = 4$ kg, $E_1 = 6.60$ N/m, $E_2 = 8.59$ N/m and $\eta = 48.9$ N·s/m. These values are used as an example and they are not calibrated to any particular valve/WDN system. The step response of the Maxwell model is plotted in Fig. 5 and shows typical effects of inertial and oscillatory components similar to those exhibited by the previously described inertial-oscillator model (see Fig. 6). The root locus plot for the closed-loop feedback control system from Fig. 1 in which the valve/WDN plant is represented by the above Maxwell model with additional actuator inertia is shown in Fig. 7. The control system has four poles marked with crosses and one zero marked with a circle. Three poles come from the Maxwell model and one from the actuator. When the gain increases from zero to infinity the dominant real pole (closest to the imaginary axes) approaches the zero, the fast actuator real pole goes to 'real' negative infinity whilst the two remaining complex poles diverge to the 'complex' infinity. The closed-system loses stability at the moment where the static gain K reaches the value of 1.41 at which point the two conjugate complex poles cross from left half plane (LHP) to right half plane (RHP). The initial value of the static gain of the Maxwell model $K_0 = 0.605$, so even a moderate increase of the gain from 0.605 to 1.41 can lead to instability and the same is likely to be true for real valve/WDN systems.

NETWORK MODEL

Model schematic

In order to simplify our theoretical static gain calculations presented in the next section as well as dynamic simulations used to recreate the real-life PRV instability event briefly described in the Introduction, the hydraulic model was reduced to just three components: the upstream pipe, the

234 valve, and the downstream pipe. The rest of the network, shown in dotted lines in Fig. 8 was not
235 modeled. Instead, it was assumed that the demands in all major nodes could be represented as
236 pressure dependent outflow from the end of the downstream pipe and modeled with a modified
237 Toricelli's orifice equation: $Q(c(t), H_4(t)) = c(t) (H_4(t) - z_4)^\alpha$ in which the pressure dependency
238 coefficient α may differ from the original Toricelli's value of 0.5 (see Ferrante et al. (2014)).
239 However it needs to be noted that the value of α in this study is likely to be different from the one in
240 Ferrante et al. (2014) since we consider the pressure at the entry to a district metering area (DMA)
241 whilst Ferrante et al. (2014) used the average pressure within the DMA. Since pressure changes at
242 the nodes far down in the network constitute just a fraction of the changes in the DMA inlet pressure,
243 the pressure dependency coefficient is likely to be lower from the ones reported by Ferrante et al.
244 (2014), although this remains yet to be proven. In the equation above z_4 denotes the elevation of
245 node 4, H_4 is the pressure head at node 4 and $c(t)$ denotes the time-varying equivalent total orifice
246 area at the end of the downstream pipe, i.e. DMA inlet (see Fig. 8).

247 The system is thus modeled as pressure-driven, not demand-driven. The demand-driven ap-
248 proach is not valid for dynamic systems because it reverses causality between flow and pressure.
249 In demand-driven approach often used in steady-state calculations the flow is forced and the nodal
250 pressures are calculated to produce the desired flows. In other words, changes in nodal pressure
251 values are the effect of flow variations, whilst in physical systems it is the flow which is driven
252 by pressure difference. Whilst such reversal of causality allows us to predict steady-state nodal
253 pressures in a network under given demands, such an approach is not valid in dynamic simulation
254 in which dynamic effects of inertia and water compressibility must not be neglected. In partic-
255 ular, demand-driven approach applied to dynamic system simulation neglects energy dissipation
256 through openings (leaks and orifices at demand nodes) which occurs when the incoming pressure
257 forces excess water outside of the system. This leads to a higher buildup of potential energy inside
258 the system, resulting in overestimation of pressure wave amplitudes – see Jung et al. (2009) for
259 comparison of both approaches on larger network models.

Model equations

Algebraic model

The algebraic steady-state model of the system depicted in Fig. 8 is composed of four algebraic equations representing pressure drops across all three components of the network: the upstream pipe, the PRV, and the downstream pipe, supplemented with flow continuity equation. The resulting system of equations can be written as $\mathbf{F}(x, \mathbf{y}) = \mathbf{0}$ where

$$\mathbf{F}(x, \mathbf{y}) = \begin{cases} H_{in} - H_u - R_1 Q^2 \\ H_u - H_d - \frac{1}{K_v(x)} Q^2 \\ H_d - H_{out} - R_2 Q^2 \\ Q - c (H_{out} - z_4)^\alpha \end{cases} \quad (5)$$

in which x denotes the valve position and $\mathbf{y} = (H_u \ H_d \ H_{out} \ Q)^T$.

Dynamic model

The dynamic simulation model of the upstream and downstream pipe is composed of two water-hammer equations solved with the well-known method of characteristics (Wylie and Streeter 1993) which transforms the two partial differential water-hammer equations into two ODEs which are solved for each internal point along the length of the pipe, i.e. $\forall i \in \{2, \dots, n-1\}$, where n is the number of nodes in the conduit. In a short form these two ODEs can be written as two characteristic equations: $C^+ : H_i = C_{p,i} - B Q_i$ and $C^- : H_i = C_{m,i} + B Q_i$. where $C_{p,i} = H_{i-1}^* + B Q_{i-1}^* - R Q_{i-1}^* |Q_{i-1}^*|$ and $C_{m,i} = H_{i+1}^* - B Q_{i+1}^* + R q_{i+1}^* |q_{i+1}^*|$ in which $B = \frac{a}{g A}$ and $R = \frac{\lambda \Delta x}{2 g D A^2}$, Δx is the distance between internal pipe nodes and * denotes the value of the variable recorded in the previous time step. The initial flow value Q_0 in the system and the piezometric head values at points 2, 3, and 4 (see Fig. 8) required for the solution of the water-hammer model are obtained from Eq. 5 for $c(t=0) = c_0$ and $x(t=0) = x_0$. All internal node flow values within the pipes are assigned the initial flow value Q_0 , i.e. $\forall i \in \{1, \dots, n\} : Q_i = Q_0$. The internal head values H_i are obtained through a linear interpolation between the boundary head values. The boundary

282 conditions for the water-hammer model are listed in Table 1. Both pipes combined are discretized
 283 into $n - 1$ segments, in which node m represents point 2, node $m + 1 =$ point 3, whilst node n
 284 represents point 4 in Fig. 8.

285 **CALCULATION OF THE STATIC GAIN OF THE VALVE/WDN SYSTEM**

286 It is argued in this paper that the static valve/WDN gain strongly depends on valve position, i.e.
 287 $K = K(x)$. The gain is low for large valve opening and increases gradually as the valve opening
 288 is reduced. The increase in the value of the gain as the valve position changes from $x = 80\%$ to
 289 $x = 30\%$ can be as high as 5 fold which presents a challenge for designing an effective controller.
 290 Since valve position influences the flow, the plant gain is also dependent on flow, i.e. $K = K(Q(x))$.
 291 As shall be shown below, for a PRV-controlled DMA under a constant H_d we can determine a
 292 unique relationship between x and Q and since Q increases as the valve is opened, K is high for low
 293 flows making the valve less stable under low-flow conditions and supporting the generally accepted
 294 opinion that PRVs become unstable under small flows.

295 From definition $K(x) = dH_d/dx$ and can be calculated from the system of equations given
 296 in Eq. 5 using the implicit function theorem (see e.g. Hubbard and Hubbard 2015) which states
 297 that if we have an implicit system of nonlinear equations $\mathbf{F}(x, \mathbf{y}) = \mathbf{0}$, such as one given in Eq. 5,
 298 and it satisfies some mild conditions on its partial derivatives, then we can calculate the vector of
 299 derivatives: $\frac{d\mathbf{y}}{dx} = -\mathbf{F}_y^{-1} \mathbf{F}_x$ where \mathbf{F}_x represents the vector of derivatives of \mathbf{F} with respect to x
 300 and \mathbf{F}_y represents the matrix of derivatives of \mathbf{F} with respect to \mathbf{y} . Through differentiation of our
 301 system of equations given in Eq. 5 we obtain:

$$302 \quad \mathbf{F}_x = \begin{pmatrix} 0 \\ 2 \frac{1}{[K_v(x)]^3} \frac{dK_v}{dx} Q^2 \\ 0 \\ 0 \end{pmatrix} \quad (6)$$

303

$$\mathbf{F}_y = \begin{pmatrix} -1 & 0 & 0 & -2 R_1 Q \\ 1 & -1 & 0 & -2 \frac{1}{[K_v(x)]^2} Q \\ 0 & 1 & -1 & -2 R_2 Q \\ 0 & 0 & -\alpha c (H_{out} - z_4)^{\alpha-1} & 1 \end{pmatrix} \quad (7)$$

304

305 where $\mathbf{y} = (H_u \ H_d \ H_{out} \ Q)^T$. By visual inspection we can see that the columns of \mathbf{F}_y are linearly
 306 independent and thus \mathbf{F}_y is invertible with a non-zero determinant:

307

$$\det \mathbf{F}_y = - \left[1 + 2 \alpha c (H_{out} - z_4)^{\alpha-1} Q \left(R_1 + R_2 + \frac{1}{[K_v(x)]^2} \right) \right] \quad (8)$$

308 We can also see that \mathbf{F}_x has only one non-zero element $\mathbf{F}_x(2)$ and thus in order to calculate
 309 dH_d/dx we only require to determine $\mathbf{F}_y^{-1}(2, 2)$. Then

310

$$\frac{dH_d}{dx} = \frac{d\mathbf{y}(2)}{dx} = -\mathbf{F}_y^{-1}(2, 2) \cdot \mathbf{F}_x(2) \quad (9)$$

311 After applying a co-factor expansion with respect to row 2 and column 2 we obtain the following
 312 relationship for $\mathbf{F}_y^{-1}(2, 2)$

313

$$\mathbf{F}_y^{-1}(2, 2) = \frac{1 + 2 R_2 Q \alpha c (H_{out} - z_4)^{\alpha-1}}{\det \mathbf{F}_y} \quad (10)$$

314 which after substitution into Eq. 9 yields the following equation for the static gain of the valve/WDN
 315 system displayed in Fig. 8

316

$$\frac{dH_d}{dx} = \frac{1 + 2 R_2 Q \alpha c (H_{out} - z_4)^{\alpha-1}}{1 + 2 \alpha c (H_{out} - z_4)^{\alpha-1} Q \left(R_1 + R_2 + \frac{1}{[K_v(x)]^2} \right)} \cdot 2 \frac{1}{[K_v(x)]^3} \frac{dK_v}{dx} Q^2 \quad (11)$$

317 where $Q = c (H_{out} - z_4)^\alpha$

318 **Static gain of an isolated PRV**

319 Before we begin to analyze Eq. 11 let us first look at a simple case of an isolated PRV where
 320 $Q = \text{const}$ and $H_u = \text{const}$ and calculate the static gain of the PRV in this isolated scenario in which

321 the PRV has no interaction with the rest of the network. We can accomplish this task two-fold.
 322 First, we can investigate a single valve equation, i.e. the second equation in Eq. 5

$$323 \quad H_u - H_d - \frac{Q^2}{[K_v(x)]^2} = 0 \quad (12)$$

324 which can be rearranged to yield an explicit relationship for downstream head, $H_d = H_u - Q^2/K_v^2$.
 325 Differentiation of H_d vs. x under the assumption that $dH_u/dx \equiv 0$ and $dQ/dx \equiv 0$ yields

$$326 \quad \frac{dH_d}{dx} = \frac{d}{dx} \left(H_u - \frac{Q^2}{[K_v(x)]^2} \right) = 2Q^2 \frac{1}{[K_v(x)]^3} \frac{dK_v(x)}{dx} \quad (13)$$

327 which shows that the valve gain is inversely proportional to the valve capacity K_v in third power.
 328 Although the term dK_v/dx has some influence on the value of the valve gain, dH_d/dx is most
 329 sensitive to K_v^3 , which shows that the change in the valve gain between low and high valve openings
 330 (e.g. nonlinearity in the control system) is not so much a result of a hydraulic nonlinearity in the
 331 valve but the fact that the valve capacity itself is low under low openings. We can also see that the
 332 isolated PRV gain described by Eq. 13 is equal to the second factor in the connected valve/WDN
 333 gain given by Eq. 11 demonstrating that the static gain of the valve/WDN system is a product
 334 of interaction between the static gain of the isolated PRV and the hydraulic characteristics of the
 335 WDN, specifically pipe resistances/conductivities and the pressure dependency coefficient of the
 336 demand(s).

337 Another way of looking at the isolated PRV gain is through the elements of the matrix \mathbf{F}_y .
 338 Under assumption that $Q = \text{const}$ the elements in rows 1, 2 and 3 in column 4 of \mathbf{F}_y become
 339 null and \mathbf{F}_y turns into a lower triangular matrix for which the determinant is a multiplication of
 340 all elements along the leading diagonal (1 in our case) and $\mathbf{F}_y^{-1}(2, 2) = 1/\mathbf{F}_y(2, 2) = -1$. Hence,
 341 $dH_d/dx = \mathbf{F}_x(2)$ and satisfies Eq. 13.

342 **Gain of the PRV connected to the network**

343 Pressure and flow in the valve/WDN system is a result of an interaction between the valve
 344 (valve characteristic) and the water distribution network (system curve). The theoretical system

345 curve depends on the assumed demand model, namely its pressure dependency characteristics. In
 346 the extreme case where demands are assumed to be forced and independent from nodal pressures
 347 the system curve $Q(H_d)$ is constant and equal to the total demand in the system. In the general
 348 case with demands depending on pressures the total flow Q through the PRV will increase with the
 349 downstream head, H_d . In such a case, i.e. with pressure-dependent demands, we can see that the
 350 static gain of the isolated PRV described by Eq. 13 and the second term in Eq. 11 is scaled by the
 351 first term in Eq. 9 which considers the properties of the WDN, i.e. pipe resistances, valve capacity
 352 and pressure dependency characteristics.

353 As shall be shown in numbers in the next section, this scaling function is always lower than
 354 unity which means that the interaction between the network and the valve leads to the reduction
 355 of the static valve/WDN gain compared to the isolated PRV gain. This behavior is very intuitive
 356 and can be easily explained. Changing valve position x leads to the change in the energy loss
 357 across the valve which in turn alters the total (potential + kinetic) energy in the system. In the
 358 isolated case, or where demands are forced, i.e. $Q = \text{const}$ the kinetic energy remains constant
 359 and hence variation in x leads to the change in potential energy only, i.e. nodal pressures. In the
 360 scenario with PRV connected to WDN and with pressure dependent demands both potential energy
 361 (nodal pressures) and kinetic energy (flow) are affected by the changes in x and thus the variation
 362 in nodal pressures will be lower as some part of the total energy is being diverted into kinetic
 363 energy. Therefore, dH_d/dx will always be lower in the valve/WDN system with pressure dependent
 364 demands compared to the isolated PRV scenario. The relationship between the closed-loop gain of
 365 the PRV connected to the network and the gain of the isolated PRV can be written as:

$$366 \quad K_{connected}(x) = f(\mathbf{p}) K_{isolated}(x) \quad (14)$$

367 where \mathbf{p} is a vector of parameters characterizing the hydraulic properties of the network, the valve,
 368 and pressure-dependency of the demands. If we look at the first term in Eq. 11 we can see that
 369 $f(\mathbf{p}) < 1$ since R_1 , R_2 and K_v are always greater than zero.

CASE STUDY

This manuscript began with a brief description of a real-life instability event which occurred in a large-scale pressure control scheme installed in a WDN of one of the major cities in the UK. This instability event shall be used here as a source of data for the case-study which aim is to test the validity of the theoretical work presented above and to showcase via simulation the applicability of the proposed remedy against instabilities for electronically controlled PRVs.

The instability event under study is shown in Fig. 2 consisting of two subplots. The top subplot shows the outlet pressure p_{out} from the PRV whilst the bottom subplot shows the valve position x . The pressure set-point set at 0.65 bar is marked with a thick dashed line and the ± 0.1 bar dead zone between which the valve element is not actuated is shown with two horizontal thick solid lines. In both subplots the time scale corresponds to the morning hours of the day where the instability event occurred. The instability began around 5:37 a.m. with relatively small pressure oscillations between 0.25 bar and 1.0 bar which then went completely out of control around 6:24 a.m. when the valve opening started to oscillate between the minimum and the maximum allowed values of 0% and 80% respectively resulting in downstream pressure variations between 0 and 4 bars. Around 6:45 a.m. the amplitude of pressure oscillations began to decrease as a result of increasing flow through the valve. Although the flow is not displayed in the plots, it can be inferred from the valve position. As shall be shown later in Fig. 10 a direct positive and mildly nonlinear relationship exists between flow and valve position in the PRV-controlled system under study and this relationship is likely to be similar for other valve/WDN systems operating under the constant outlet pressure setpoint. Around 8:20 a.m. the oscillations decayed and the system resumed normal operation and began to keep the outlet pressure at the setpoint value of 0.65 ± 0.1 bar. This return to stability coincided with an increase in flow rate through the valve.

The PRV under study is fed by a fixed head reservoir with static pressure head $H_{in} = 186.5$ m connected to the PRV through ~ 10 km pipe of 0.8 m diameter. The PRV is set to drop the upstream pressure to 106.5 m which, considering pressure losses in the upstream pipe, results in an average pressure head drop of about 75 m. The downstream pipe goes for a distance of about 2

397 km after which it starts dividing into a dense network of pipes feeding different DMAs in the city.
398 The PRV is controlled electronically by a PLC implementing a time-discrete proportional integral
399 derivative controller (PID).

400 The work presented in this section proceeds as follows. First, pipe resistances of the simplified
401 static pipe-PRV-pipe model from Fig. 8 are calibrated, given the PRV characteristic obtained from
402 the measurements, such that the static hydraulic characteristic of the model matches that of the real
403 network. This calibrated simplified static model is then used to derive the static gain curves of the
404 valve/WDN system and to calculate the formula for the gain compensator which aim is to keep the
405 static gain of the system constant within the entire permissible valve opening range. The positive
406 effects of the gain compensator on the valve stability are then tested via dynamic simulation.

407 **Valve capacity curve**

408 The valve capacity curve was provided by the manufacturer and additionally calculated from 15-
409 min measurements of flow, upstream and downstream pressures, and valve position. As pictured
410 in Fig. 9 the manufacturer's capacity curve and the measured capacity curve are significantly
411 different indicating that either the valve was not made to the specification or, more likely, that
412 the valve's capacity had changed over time due to wear and tear of internal components. The
413 manufacturer's capacity curve and the capacity curve from the measurements were obtained via
414 nonlinear curve fitting and are plotted in Fig. 9 with a thick solid and thick dash-dotted line,
415 respectively. The obtained capacity curve equations are as follows. The manufacturer's valve
416 capacity curve: $K_{v,manu} = 0.2360 x^2 + 0.2491 x$, the valve capacity curve from the measurements:
417 $K_{v,meas} = 0.1597 x^2 - 0.01129 x$ The shaded surface between both curves represents the valve
418 capacity uncertainty region indicating that the change of valve capacity over time needs to be
419 accounted for at a controller design stage.

420 **Relationship between valve position and flow**

421 Data points obtained from 15-min measurements of flow and valve position and plotted in Fig. 10
422 as gray dots show that for our valve/WDN system a direct and slightly nonlinear relationship exists
423 between valve position and flow. This relationship was found to be quadratic with the line of best fit

424 given as $Q(x) = 0.1155 x^2 + 0.3458 x$ which is marked in Fig. 10 with a thin dashed line. From the
425 control point of view the existence of such a relationship allows us to infer the flow rate from the
426 valve position or vice versa. This additional information can be used as a surrogate measurement
427 for control purposes or sensor fault diagnosis. In this paper, valve position and flow measurements
428 obtained from the case study were used to calibrate the static hydraulic properties of our simplified
429 pipe-PRV-pipe model shown in Fig. 8 such that the model and the real network exhibit similar
430 hydraulic characteristics. For this purpose we used the valve capacity curve obtained from the
431 measurements, as explained in the section above. Our simplified model is governed by the system
432 of equations given in Eq. 5. As we can see, the parameters which can be used to calibrate the
433 hydraulic characteristics are the upstream and downstream pipe resistances R_1 , R_2 and the pressure
434 dependency coefficient α . During calibration it was found that α was not a sensitive parameter.
435 Hence, for simplicity, we assumed $\alpha = 0.5$, as in the Toricelli's equation.

436 Since the pipe resistance coefficients, i.e. R_1 and R_2 lump three parameters characterizing the
437 pipes' hydraulics, i.e. length L , absolute roughness ϵ and diameter D we cannot uniquely find the
438 right combination of these three parameters in this study. They are nevertheless very important
439 in dynamic simulations as they affect in different ways how the pressure builds up and propagates
440 across the network. The product of L and D^2 (multiplied by $\pi/4$) determines the total volume
441 of water in the pipes and thus the amount of inertia in the system. The higher the inertia the
442 higher the amplitude of pressure waves produced during valve closure/opening as the forces acting
443 on the closing element are proportional to the accelerating/decelerating mass. The pipe length L
444 determines the period of pressure waves departing from/arriving at the PRV for a given wave speed.
445 The absolute pipe roughness ϵ determines the amount of dissipative friction forces in the system
446 and affects the pressure drop in the pipe in steady-state as well as the amount of pressure wave
447 attenuation during transient flow. Bearing this in mind the calibrated parameters R_1 and R_2 were
448 first adjusted such that the model characteristics matched the measurements. Then the information
449 about pipe friction coefficients and upstream and downstream pipe diameters was taken from site,
450 allowing the authors to adjust the pipe lengths L_1 and L_2 . In case some of the above information is

451 not available and need to be assumed the results of the transient simulations may differ significantly
452 from the measurements in terms of pressure wave period, pressure wave amplitudes and amplitude
453 attenuation. These discrepancies may be reduced during further calibration on a dynamic model by
454 adjusting, specifically, the values of pipe lengths, pipe diameters, and pipe roughness coefficients.

455 As mentioned above, the pressure dependency coefficient α was found to be an insensitive
456 parameter. The verbal explanation of this fact is as follows. An increase in α will produce a higher
457 flow for a given value of pressure head H_4 or orifice opening c (see Fig. 8). If the flow increases and
458 H_3 is kept constant, i.e. by the PRV, then the valve needs to open in order to reduce the pressure
459 drop which is proportional to Q^2 . Hence Q and x will increase/decrease in the same direction and
460 the Q vs. x curve will not change by much. Nevertheless, we should bear in mind that the value of
461 α has a significant impact on the transient flow as it determines how much potential energy entering
462 with the wave as pressure is lost in the system as the water is ejected through the orifices (openings)
463 at individual nodes or, in our case, at the end of the downstream pipe.

464 The effects of R_1 and R_2 on the Q vs x characteristic of our simplified hydraulic model are
465 as follows. If R_1 is increased then the pressure drop in the upstream pipe will be more sensitive
466 to flow, thus H_2 will decrease more with respect to Q and hence we will need to have a higher
467 valve opening x for a given flow, since the pressure drop across the PRV will need to be smaller
468 to maintain constant downstream pressure. The curve will move left if R_1 is increased. If R_2 is
469 increased then the pressure drop in the downstream pipe will be more sensitive to flow, thus H_4
470 will be lower for a given flow and the flow, since it's pressure dependent, will decrease. Hence, we
471 will have lower flows for the same opening. The curve will move down if R_2 is increased.

472 The calibration was performed manually. The obtained model properties are as follows: $L_1 =$
473 $5,000$ m, $L_2 = 10,000$ m, $D_1 = D_2 = 0.8$ m, $\epsilon_1 = \epsilon_2 = 0.003$ m, $z_4 = 50$ m. The calibrated
474 curve is shown in Fig. 10 with a thick dash-dotted line. As we can see, the curve matches the
475 measurements and is closely aligned with the line of best fit (thin dashed line). Additionally, we
476 also plotted the curve for $K_v = K_{v,manu}$, i.e. for the manufacturer's valve capacity characteristic.
477 Since the manufacturer's valve capacity is higher from the measured one the system would allow

478 higher flows for a given valve opening and therefore the curve (thick solid line) lies above the
479 measurements.

480 **Valve gain analysis**

481 The calibrated parameter values and the measured valve characteristic were fed into Eq. 11 and
482 used for subsequent calculation of the static valve/WDN gain vs. valve opening for different values
483 of α . Additionally, Eq. 13 was used to calculate the isolated PRV gain, i.e. for $Q = \text{const}$ and
484 $H_u = \text{const}$. The resulting curves were plotted in Fig. 11. As predicted, the static valve/WDN gain
485 decreases with valve opening. The difference in gain between low and high valve openings is most
486 severe for the isolated PRV case and becomes smaller in the connected valve/WDN case due to the
487 reducing effects of the WDN, i.e. dependency of flow on pressure. This interaction between the
488 PRV and the rest of the network was already explained earlier on in the text - see Eq. 14. Fig. 11
489 shows that the gain vs. valve position curve becomes less steep as α increases. Whilst the gain can
490 change from 1 at $x = 80\%$ to 6 at $x = 20\%$ for $\alpha = 0.5$, this change can be as severe as 1 to 10 for
491 $\alpha = 0.3$.

492 The gain values were calculated for two valve characteristics: the manufacturer's characteristic
493 $K_{v,manu}$ (shown in solid line) and the measured characteristic $K_{v,meas}$ (shown in dashed line). As
494 can be seen, both curves lie close together for medium valve openings between 30% and 60% and
495 began to diverge at extreme values close to the minimum and maximum allowed openings of 20%
496 and 80% respectively. The valve capacity has a larger impact on the Q vs. x characteristic (see
497 Fig. 10) which means that the flows will be different for the same valve openings. Specifically,
498 given that both PRV gain characteristics are similar and the gain is such that the valve begins to
499 oscillate at around 30%, the valve with the manufacturer's characteristic (e.g. brand new valve) will
500 begin to oscillate at a higher flow than the valve with the measured characteristic (e.g. old valve).

501 The gain curve for $K_v = K_{v,meas}$ and $\alpha = 0.5$ was approximated with a third order polynomial
502 using nonlinear curve fitting. The obtained gain equation is given below.

$$503 \quad K_{connected}(x) = K(x) = -1.201 \times 10^{-5} x^3 + 3.162 \times 10^{-3} x^2 - 0.3186 x + 12.23 \quad (15)$$

504 Eq. 15 shall be later used to formulate a nonlinear gain compensator used to prevent PRV instability
 505 under low flows conditions.

506 Fig. 12 shows the impact of our simplified network model parameters, i.e. R_1 , R_2 and the valve
 507 characteristic $K_v(x)$ on the scaling function $f(\mathbf{p})$ featured in Eq. 14. For this purpose the first term
 508 of Eq. 11 representing $f(\mathbf{p})$ was split into three different individual terms: $f_{R_1} = 2R_1Q dQ/d\alpha$,
 509 $f_{R_2} = 2R_2Q dQ/d\alpha$, $f_{K_v} = 2(1/K_v(x))^2Q dQ/d\alpha$ where $dQ/d\alpha = \alpha c (H_{out} - z_4)^{\alpha-1}$. Using these
 510 three new terms $f(\mathbf{p})$ can now be represented as:

$$511 \quad f(\mathbf{p}) = f(R_1, R_2, K_v(x), \alpha) = \frac{1 + f_{R_2}}{1 + f_{R_1} + f_{R_2} + f_{K_v}} \quad (16)$$

512 f_{R_1} , f_{R_2} and f_{K_v} were plotted in Fig. 12 for $\alpha = 0.3$ and $\alpha = 0.5$. As shown, all three terms show a
 513 stronger relationship vs. x for higher α values. We can also see that the most dominant parameter
 514 is f_{K_v} , although the scaling effects of R_1 and R_2 become significant at higher openings. What this
 515 means is that if, i.e. R_1 or R_2 are increased the gain value at lower openings will remain similar
 516 whilst the gain reduction at larger openings shall be higher. The overall scaling factor f for $\alpha = 0.3$
 517 and $\alpha = 0.5$ was plotted in thick solid and thick dashed line respectively. As explained earlier using
 518 Eq. 14 the scaling factor is always less than unity.

519 **Control loop structure**

520 The valve/WDN control loop structure (see Fig. 13) and controller parameters are as follows.
 521 The valve with a known capacity characteristic described above is controlled with a time discrete
 522 PID with the following gains: $K_p = 0.5\%/m$, $K_i = 0.05\%/(m\ s)$ (i.e. $T_i = 10s$), $K_d = 0(m\ s)/\%$;
 523 and with sampling time $t_s = 0.1s$. The actuator is modeled as a first order system with time constant
 524 of $0.1s$. The valve opening and closing rates are limited at $\pm 100/87\ \%/s$ to reduce the risk of
 525 transients. The rate-limited actuator output then goes through a backlash block with a deadband of
 526 0.8% . The rate limiter and the deadband are embedded within the Actuator block and are not shown
 527 in the block diagram. The controller error e is the difference between the downstream pressure
 528 head setpoint $H_d^{set} = 106.5m$ and the downstream pressure measurement $H_{d,ave}^{meas}$ (m) averaged using

529 a moving average filter with a buffer size of 300 data points updated every 0.02s and followed by
 530 a zero-order-hold of 0.1s. To reduce the amount of control effort the control system is only active
 531 when $|e| \geq 0.5\text{m}$, i.e. the outlet pressure head H_d is allowed to vary $\pm 0.5\text{m}$ - see Dead zone block
 532 in Fig. 13.

533 **Nonlinear compensator for stability improvement**

534 One way to combat the nonlinearity in the valve/WDN system resulting from the dependence
 535 of static gain on valve position is to introduce a nonlinear compensator in the forward path of
 536 the feedback system which will scale the control signal depending on the current measurement
 537 of the valve position, such that the gain is constant (or near constant in practice) over the entire
 538 permitted range of valve openings. Placement of the nonlinear compensator in the feedback system
 539 structure is indicated in Fig. 13 with a dashed rectangle. The compensator receives the valve
 540 position measurement from the valve actuator and calculates the correcting factor either from a
 541 look-up table or from an analytical formula, such as the one derived below for our system. It
 542 is recommended that the PID is tuned at the typical valve position x_{typ} around 40%-50%. The
 543 correcting factor can then be calculated from the following formula.

$$544 \quad k(x) = \frac{K_p(x = x_{typ})}{K_p(x_{meas})} \quad (17)$$

545 where x is the actual valve position, x_{meas} is the measured valve position and x_{typ} is the typical
 546 valve position for which the controller was tuned. If during operation of the control system the
 547 valve position changes from x_{typ} to x , the gain, as seen by the controller, will theoretically remain
 548 the same as for $x = x_{typ}$.

$$549 \quad K = K(x) k(x) = K(x) \frac{K(x = x_{typ})}{K(x_{meas})} \approx K(x = x_{typ}) \quad (18)$$

550 The above equation is valid if $x_{meas} \approx x$, i.e. we have an accurate measurement of valve position and
 551 we have a perfect compensation model $k(x)$ under all operating conditions. In practice, $K \neq \text{const}$
 552 since $x \neq x_{meas}$ and the compensator model will never be perfect. However, in theory, assuming

553 perfect valve position measurements and perfect gain compensation the PID gains, i.e. K_p , K_i and
 554 K_d tuned for the typical valve position should, in theory, be valid and satisfying the control system
 555 design criteria across the entire operating range of the PRV. The static gain correction formula
 556 was obtained from Eq. 17 assuming the model hydraulic structure in Fig. 8, $\alpha = 0.5$, $K_v = K_{v,meas}$
 557 and the typical valve position $x_{typ} = 50\%$ for which the fitted polynomial curve (see Fig. 11) was
 558 calculated to be:

$$559 \quad k(x) = \frac{K(x_{typ} = 50\%)}{K(x)} = \frac{2.340}{-8.280 \times 10^{-6} x^3 + 2.450 \times 10^{-3} x^2 - 0.2658 x + 10.54} \quad (19)$$

560 $k(x)$ curve was plotted in Fig. 14. The plot shows that the scaling function $k(x) \equiv 1$ for $x = x_{typ}$,
 561 in our case $x = 50\%$ whilst $k(x) < 1$ for $x < x_{typ}$ and $k(x) > 1$ for $x > x_{typ}$, thus maintaining the
 562 plant gain at a constant value equal to $K(x = x_{typ})$.

563 **Valve/WDN simulation with and without static gain compensator**

564 The purpose of the transient simulation described in this section is to recreate the instability
 565 event visualized in Fig. 2, and to show the improvement in the stability of the control system after
 566 the introduction of the static gain compensator. We aim to demonstrate the validity of the theoretical
 567 concepts introduced earlier in this manuscript and derived via static analysis by testing them in a
 568 transient simulation scenario with control, actuation and sensing loop taken from the real-life case
 569 study. Although the transient model was not properly calibrated on dynamic data we shall see that
 570 the simulation results and the actual instability event are indeed very similar. The second aim of
 571 this simulation study was to demonstrate that instabilities in valve/WDN systems can occur with
 572 zero disturbance, i.e. with smooth inputs, in order to reinforce the idea that the loss of stability in
 573 PRVs is mainly a result of changing static plant gain and is not in any way caused by transients,
 574 although an incoming pressure wave or change in demand can trigger the instability earlier if the
 575 valve/WDN system is already operating close to instability. A similar simulation study was already
 576 performed by Janus and Ulanicki (2017).

577 The simulations were performed on the simplified hydraulic model shown schematically in

578 Fig. 8 using water-hammer equations solved with the method of characteristics, given the boundary
579 conditions in Table 1 and assuming that the wave speed $a = 1,200$ m/s. The hydraulic parameters,
580 i.e. $L_1, L_2, D_1, D_2, \epsilon_1$ and ϵ_2 were calibrated earlier from the information provided in the measured
581 Q vs. x curve given in Fig. 10 and from site information about pipe roughness and diameters. The
582 model was created and simulated in Simulink™ and MATLAB 2015a™.

583 The control, actuation and sensing loop shown in Fig. 13 was modeled as follows. The dead-
584 zone was described with a User Specific Function block with code $(\text{abs}(u) > \text{dead_zone}) * u$ where
585 u is the block input. The rate limit on the actuator is implemented as a Rate Limiter block whilst
586 the actuator backlash is implemented using the Backlash block. The discrete PID controller was
587 modeled using the PID Controller Block with Backward Euler integration method and minimum
588 and maximum output limits at 10% and 80% respectively. The pipes were discretized such that an
589 appropriate ratio between the selected time step $\Delta t = 0.02$ s and the spacial step Δx was maintained
590 allowing proper capture of the characteristics at internal pipe nodes given the wave speed a .

591 The instability, with no gain compensator, was recreated by reducing the open orifice area at the
592 end of the downstream pipe in a linear fashion over 1 hr from an initial value $A_{init} = 1.3141 \times 10^{-2}$ m²
593 to the final value of 3.141×10^{-3} m², i.e. by 0.01 m². $A_{orifice}$ was then held at the final value for
594 the next 30 minutes after which the ramp reversed direction and the orifice area increased back to
595 the initial value in the next hour. Hence, the total simulation time was equal to 2.5 hrs but in order
596 to better display the occurring oscillations the limits on the x axis in the simulation results plots
597 (see Fig 15) were reduced to 0.5 h and 2.0 h respectively.

598 The top subplot in Fig. 15 shows that as the orifice area gradually got smaller so was the total
599 flow q_{PRV} in the network. In order to maintain a constant downstream pressure head (H_3) whilst the
600 flow was being reduced the valve was closing in an almost linear fashion. In the scenario without
601 the gain compensator (dark solid line) the pressure was within the allowed band for the first hour
602 until the valve opening reached the value of about 30% under which the PRV lost stability and began
603 to oscillate. The oscillations were rising in amplitude for the next half an hour whilst the open
604 orifice area was kept at a constant and minimum value. The valve began to slowly regain stability,

605 which manifested itself in a gradual reduction of the oscillation amplitude as $A_{orifice}$ was being
606 increased and finally resumed normal operation at $t \sim 1.8$ h when the valve opening increased to
607 about 40%. This simulated valve behavior is very similar to the observed real-life instability event
608 (see Fig. 2) where the valve lost stability at $x = 30\%$ and regained stability at $x = 60\%$.

609 In the case with the static gain compensator in place (light solid line) we can see that the valve
610 did not oscillate at low flow since the gain compensation formula lowered the static gain such that
611 it didn't reach the critical value for which the system becomes unstable. However, it has to be noted
612 that this gain compensation task was achieved with a perfect compensator, i.e. without a mismatch
613 between the simulation model and the model used to derive the compensator. Additionally, a very
614 simple hydraulic model was used in the study which represented all distributed nodal demands
615 by one time-varying orifice area. In real-life applications care must be taken that the uncertainty
616 with regards to model structure and model parameters is taken into account during the design of
617 a PRV controller or otherwise an appropriate robust field tuning method needs to be developed.
618 Even though we had a perfect compensation model we can see in the middle subfigure in Fig. 15
619 that the gain compensated valve/WDN system also begins to oscillate, albeit very slightly, between
620 1 and 1.1 hrs. It is suspected that these small oscillations are due to a limit cycle produced by
621 the nonlinearities in the control-loop, i.e. dead-zone, backlash, and 0.1s zero-order-hold. It is
622 suspected that the most significant nonlinearity at small openings is due to the backlash which,
623 having a fixed value of 0.8%, becomes proportionally large when the opening is small. Also, the
624 ± 0.5 m dead-zone might have contributed to pressure variations, since it is a well-known fact that
625 dead-zones in feedback control loops lead to limit cycles.

626 Based on the results of the above case study and the theoretical developments described in
627 the first half of this manuscript we can formulate the following practical recommendations for a
628 hydraulic engineer working in pressure management of WDNs. First, with all pressure control
629 schemes incorporating a valve of any type, an engineer should be aware that the static gain between
630 valve opening and outlet pressure can increase significantly as the valve opening is reduced, which
631 in turn can lead to valve/WDN system instability under low flow conditions. The engineer should

632 therefore first assess the gain at the nominal valve opening via theoretical calculation or from
633 online measurements and later check the gain at the minimum flow conditions. If the difference
634 between the nominal and the maximum gain is large, one should perform further analysis of
635 system stability. The gain can be calculated from the valve characteristic curve combined with
636 the network characteristics. However since the latter is difficult to obtain it is recommended that
637 downstream pressure responses to step change e.g. +0.5% in valve opening are recorded in open-
638 loop configuration, i.e. with disconnected feedback loop, at different flow rates. By taking the
639 ratio between the pressure increase and the valve position change $\Delta H_d/\Delta x$ in steady-state we will
640 calculate the static gain vs. valve position curve which can be used to design a gain compensator to
641 be used in line with an electronic controller such that the overall static gain of the system remains at
642 approximately constant value over an entire flow range. For hydraulic PRVs in which the controller
643 is implemented in hardware it is necessary to introduce modifications to the pilot control loop in
644 order to achieve similar compensation effects. This is the subject of the authors' current research.
645 In cases where a modification to the control loop cannot be made, the risk of instability can be
646 reduced by narrowing down the operating range of the control valve, such that the opening does not
647 exceed the 30% – 70% range. If due to high variability of flows the range of openings is high it may
648 be possible to install a cascade of valves in which the upstream valve is a normal actuated valve
649 with a time-scheduled opening trajectory following a diurnal demand pattern whilst the downstream
650 PRV acts as a second step in pressure reduction whilst maintaining a constant outlet pressure.

651 **CONCLUSION**

652 This manuscript demonstrated that the static gain of a valve/WDN system decreases with valve
653 opening, making outlet pressure more sensitive to valve position changes at low openings. This
654 increase in valve gain at small valve openings and thus, flows, can alone explain why valves tend to
655 become unstable under low-flow conditions, which was proved through simulation. The changing
656 nature of the static valve/WDN gain is due to an inherent property of a valve whose capacity
657 increases with valve opening, as well as hydraulic properties of the rest of the water distribution
658 network. For simple networks an analytical relationship between the static valve/WDN gain and

659 valve opening can be derived, whilst for larger networks an empirical procedure will need to be
660 developed in order to be able to measure the static gain variability and then use this information
661 to design countermeasures preventing the PRVs from oscillating. One such countermeasure,
662 i.e. static gain compensator was tested here and proved to eliminate valve instability under low
663 flow conditions. Such gain compensation is an easy solution for electronically controlled valves,
664 however care must be taken that a robust solution to such compensation formula is sought in
665 the future which is resilient to biased, delayed and noisy valve position measurements and to the
666 mismatch between the network model and the real system. Similar compensation solution can be
667 sought for hydraulically controlled PRVs, although the problem is likely to be more complicated
668 as the control system there is realized mechanically which is harder to modify than programmable
669 electronic controllers. It is envisaged that through appropriate alterations to control structures of
670 both types of valves we can make them more stable and hence reduce the risk of occurrence of
671 PRV-related and destructive transient effects in water distribution networks. It is possible that
672 further improvements can be made through appropriate selection of the type of the control valve
673 and its sizing in order to reduce the nonlinearity and narrow down the operating range of the valve
674 hence reducing the difference between the minimum and the maximum static gain. However, the
675 problem of varying static gain will still persist and will have to be accounted for during controller
676 design, e.g. via static gain compensation.

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679 for the Case Study.

680 **NOTATION**

The following symbols are used in this paper:

	A	=	pipe cross section area;
	$A_{orifice}$	=	orifice opening area;
	a	=	pressure wave speed;
	c	=	equivalent total orifice area at the end of the downstream pipe;
	D	=	pipe diameter;
	e	=	error signal;
	f	=	gain scaling factor;
	$f_{R_1}, f_{R_2}, f_{K_v}$	=	gain correcting factors due to upstream pipe, downstream pipe and valve capacity;
	G	=	transfer function;
	H	=	pressure head;
	$H_u, H_d, H_d^{set}, H_{out}, H_{in}$	=	upstream, downstream, downstream setpoint, outlet, and inlet pressure head;
	K	=	plant gain;
	$K_{connected}, K_{isolated}$	=	connected and isolated PRV gain
	K_e	=	fluid's bulk modulus of elasticity;
	K_p, K_i, K_d	=	proportional, integral, and derivative gains of a PID controller;
	K_v	=	valve capacity;
	$K_{v,manu}, K_{v,meas}$	=	valve capacity from manufacturer's data and measurements;
	k	=	gain compensation coefficient;
	L	=	pipe length;
682	p_{out}	=	outlet pressure;
	Q	=	flow;
	R	=	pipe resistance;
	s	=	Laplace variable;
	T_i	=	PID controller's integral time constant;
	t	=	time;
	t_d	=	time delay;
	t_s	=	sampling time;
	u	=	control signal;
	$x, x_{typ}, x_{meas}, x_0$	=	valve position, typical valve position, measured valve position, initial valve position;
	y	=	output;
	z	=	elevation;
	α	=	pressure dependency coefficient;
	ϵ	=	absolute pipe roughness;
	λ	=	Darcy-Weisbach friction coefficient;
	ω_n	=	natural frequency;
	ρ	=	fluid density;
	τ	=	time constant;
	ζ	=	damping ratio;

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740

List of Tables

741

1 Boundary conditions for the water-hammer model. 33

TABLE 1. Boundary conditions for the water-hammer model.

Upstream reservoir	PRV upstream
$H_1 = H_{tank}$ $Q_1 = \frac{H_{tank} - C_{m,1}}{B_1}$	$H_m = C_{p,m} - B_1 Q_m$ $Q_m = 0.5 K_v^{-2} \left(-(B_1 + B_2) + \sqrt{(B_1 + B_2)^2 - 4 K_v^{-2} (C_{m,m+1} - C_{p,m})} \right)$
PRV downstream	Outlet orifice
$H_{m+1} = C_{m,m+1} + B_2 Q_{m+1}$ $Q_{m+1} = Q_m$	$H_n = C_{p,n} - B_2 Q_n$ $Q_n = g c^2 \left(-B_2 + \sqrt{B_2^2 - \frac{2}{g c^2} (z_n - C_{p,n})} \right)$

List of Figures

743	1	Block diagram representation of a closed-loop feedback system including valve/WDN plant and a controller.	35
744			
745	2	PRV instability event recorded in a large-scale pressure control scheme in one of the major cities in the UK.	36
746			
747	3	Block diagram of the lumped inertial-oscillatory model of a valve/WDN system. . .	37
748	4	Schematic diagram of the physical Maxwell model.	38
749	5	Step response of the Maxwell model to a unit step in the input.	38
750	6	Step response of lumped inertial-oscillator model of a valve/WDN to a unit step in the input.	39
751			
752	7	Root-locus of the closed-loop system with the Maxwell model plus actuator inertia. .	40
753	8	Schematic of the hydraulic model used for dynamic simulation and calculation of static gain.	41
754			
755	9	Valve capacity according to the manufacturer's data and the measurements. Dashed lines denote 95% confidence intervals for the fitted curves.	42
756			
757	10	Relationship between total flow Q and valve position x for the distribution system under study and the simplified pipe-PRV-pipe model.	43
758			
759	11	Valve/WDN gain K against valve position x for different values of pressure dependency coefficient α	44
760			
761	12	Impact of R_1 , R_2 and K_v on the scaling function $f(\mathbf{p})$ between $K_{connected}(x)$ and $K_{isolated}(x)$	45
762			
763	13	Valve/WDN closed-loop model structure adopted from real-life system and implemented in the mathematical model used for the transient simulation study.	46
764			
765	14	Nonlinear gain compensator curve $k(x)$ used for scaling the Valve/WDN plant gain K in the transient simulation study.	47
766			
767	15	Results of the transient simulations of our simplified pipe-PRV-pipe hydraulic model with (red) and without (black) gain compensation.	48
768			

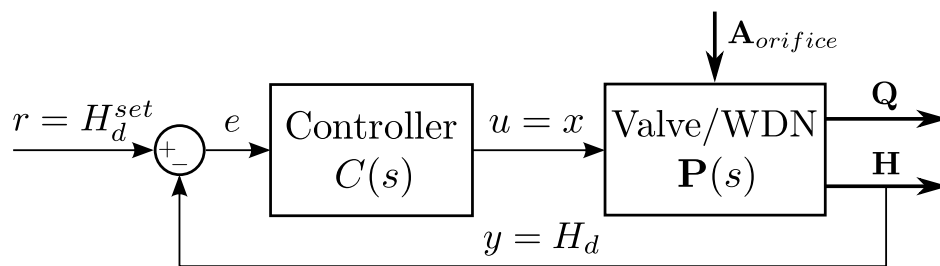


Fig. 1. Block diagram representation of a closed-loop feedback system including valve/WDN plant and a controller.

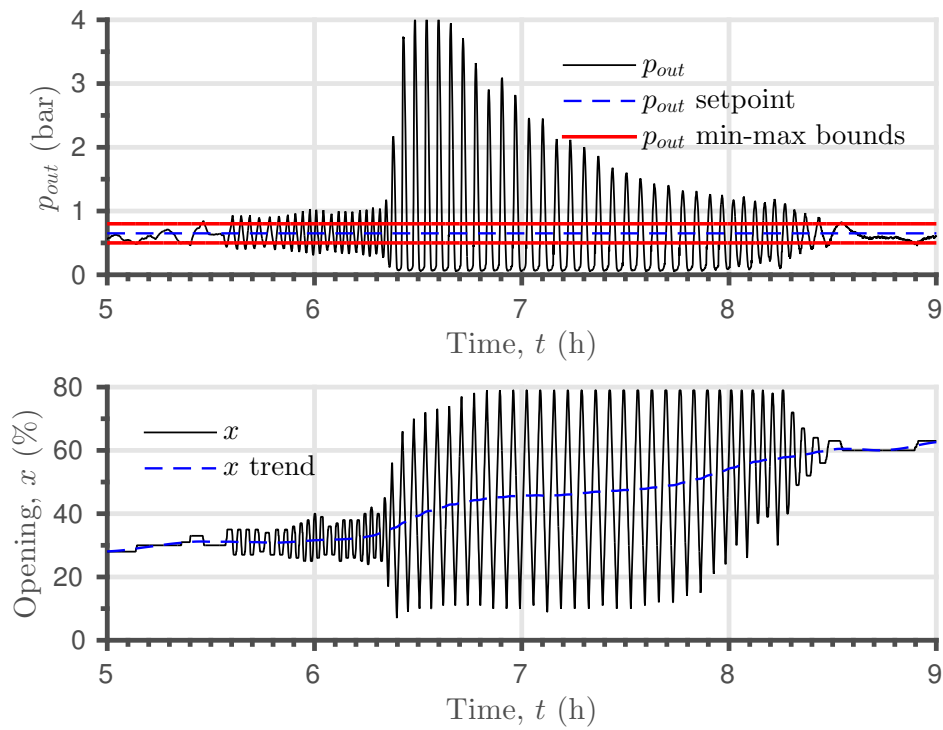


Fig. 2. PRV instability event recorded in a large-scale pressure control scheme in one of the major cities in the UK.

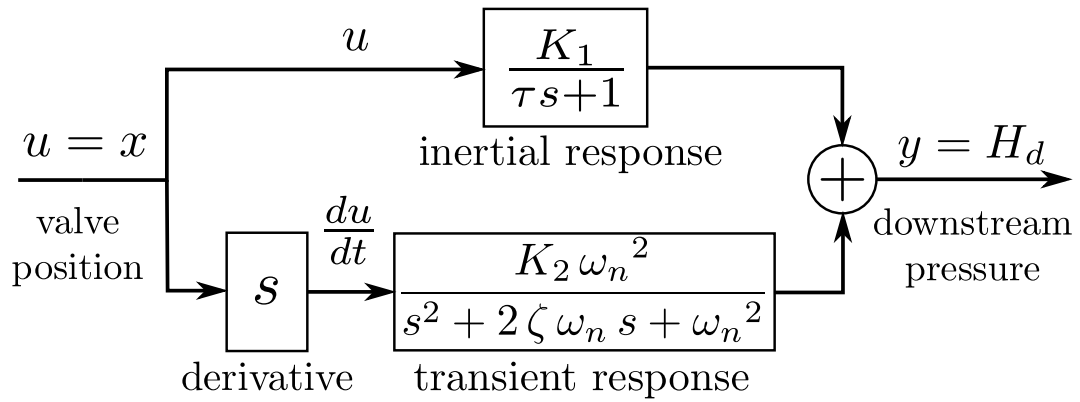


Fig. 3. Block diagram of the lumped inertial-oscillatory model of a valve/WDN system.

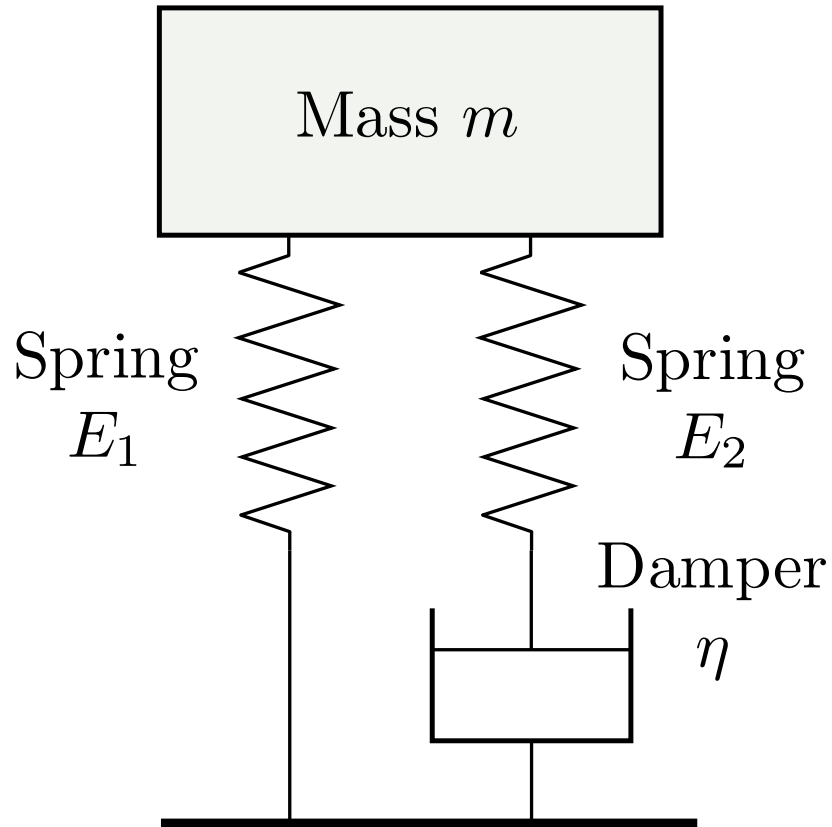


Fig. 4. Schematic diagram of the physical Maxwell model.

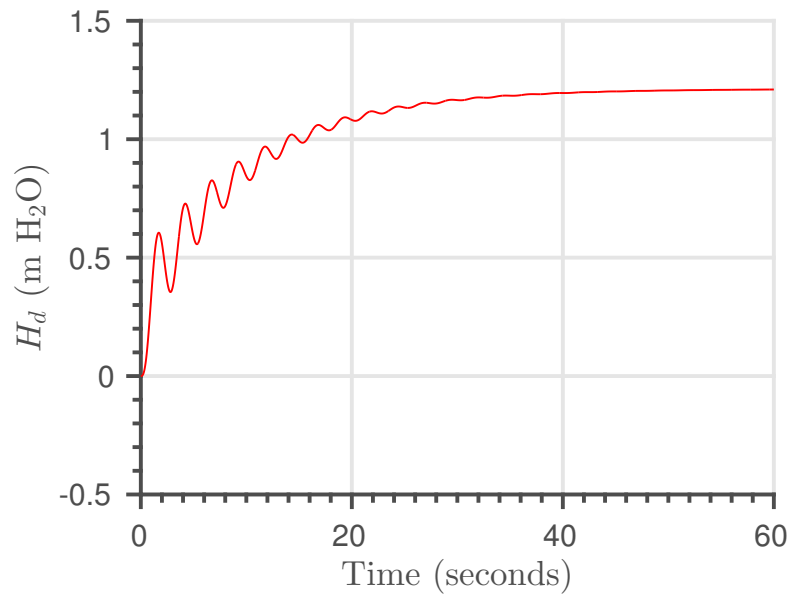


Fig. 5. Step response of the Maxwell model to a unit step in the input.

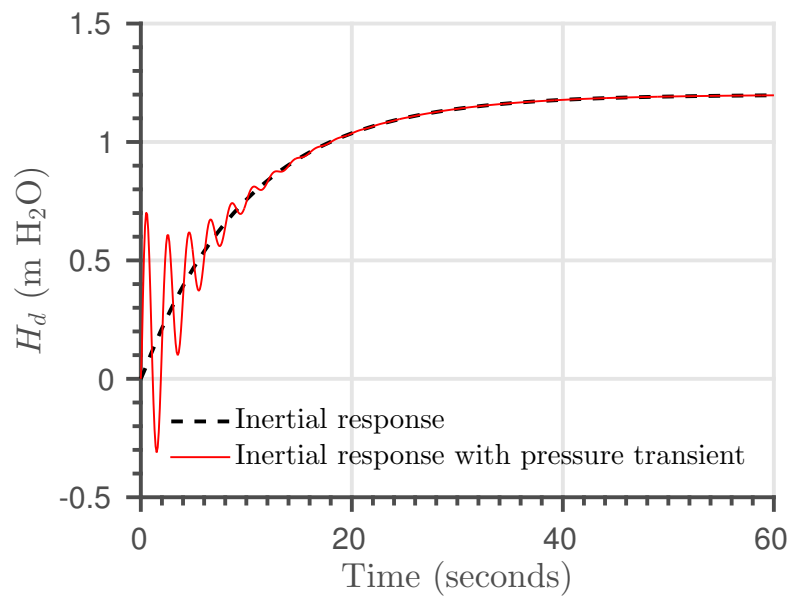


Fig. 6. Step response of lumped inertial-oscillator model of a valve/WDN to a unit step in the input.

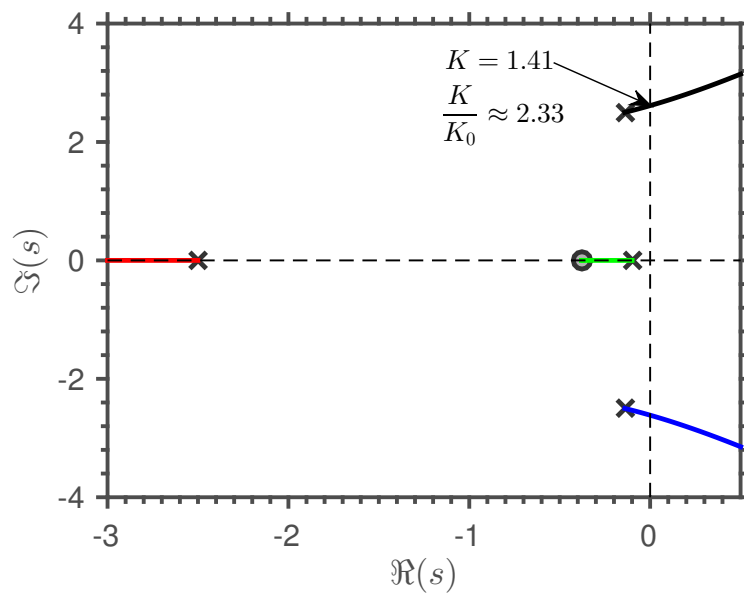


Fig. 7. Root-locus of the closed-loop system with the Maxwell model plus actuator inertia.

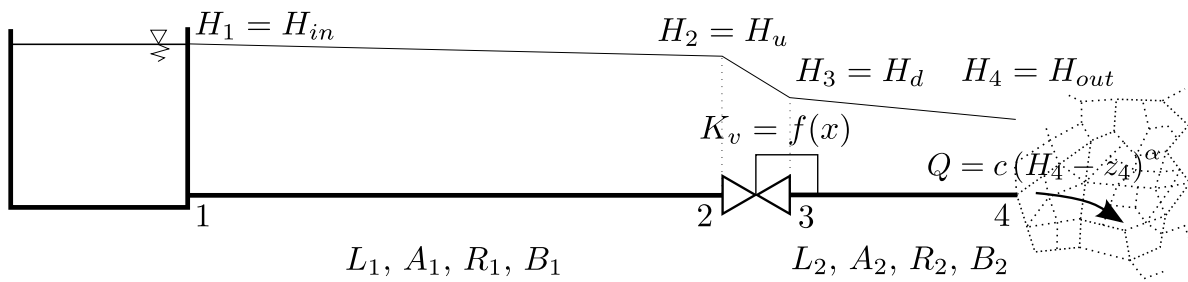


Fig. 8. Schematic of the hydraulic model used for dynamic simulation and calculation of static gain.

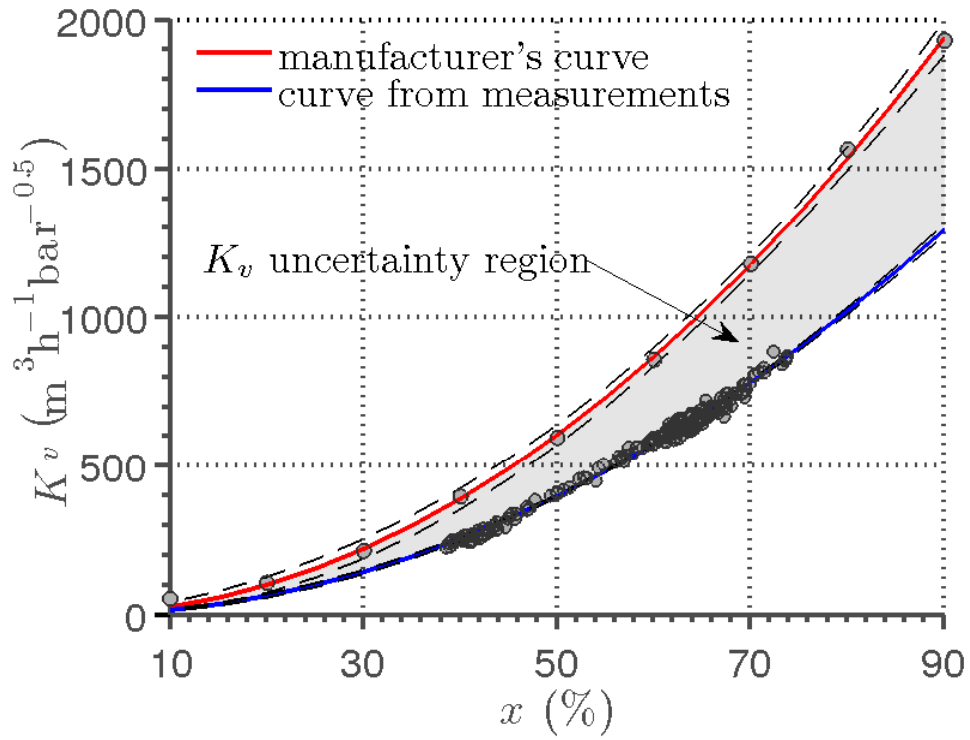


Fig. 9. Valve capacity according to the manufacturer's data and the measurements. Dashed lines denote 95% confidence intervals for the fitted curves.

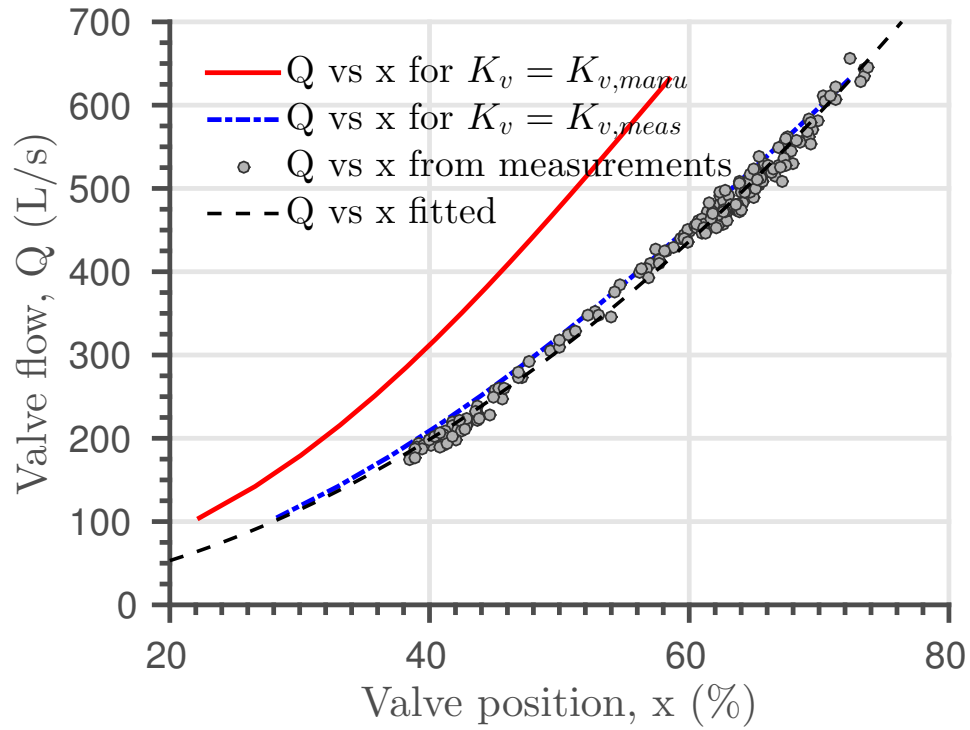


Fig. 10. Relationship between total flow Q and valve position x for the distribution system under study and the simplified pipe-PRV-pipe model.

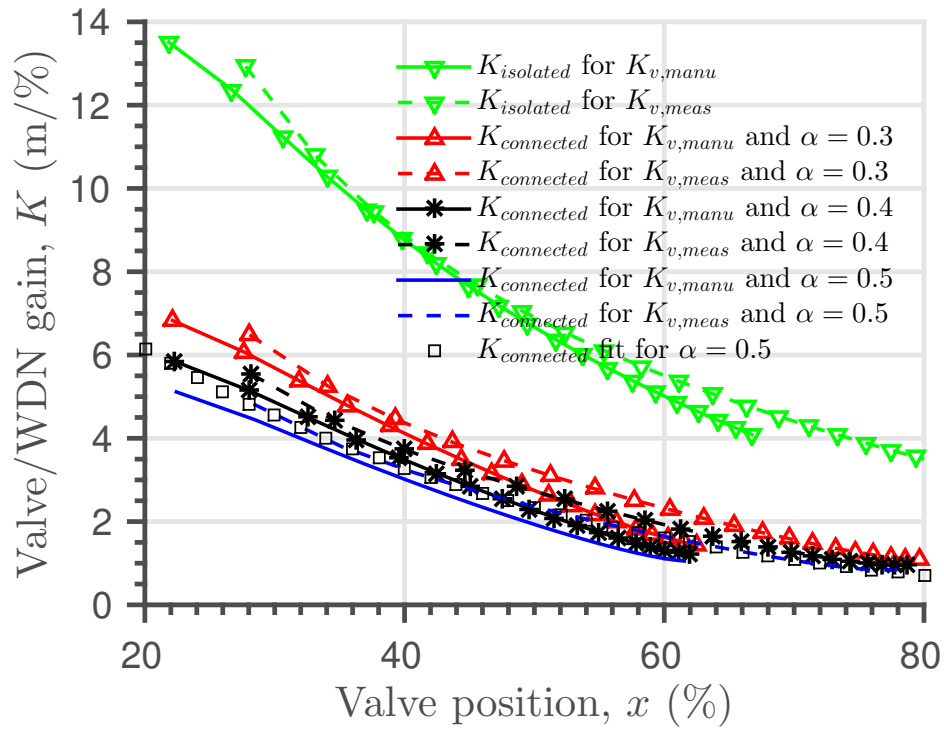


Fig. 11. Valve/WDN gain K against valve position x for different values of pressure dependency coefficient α .

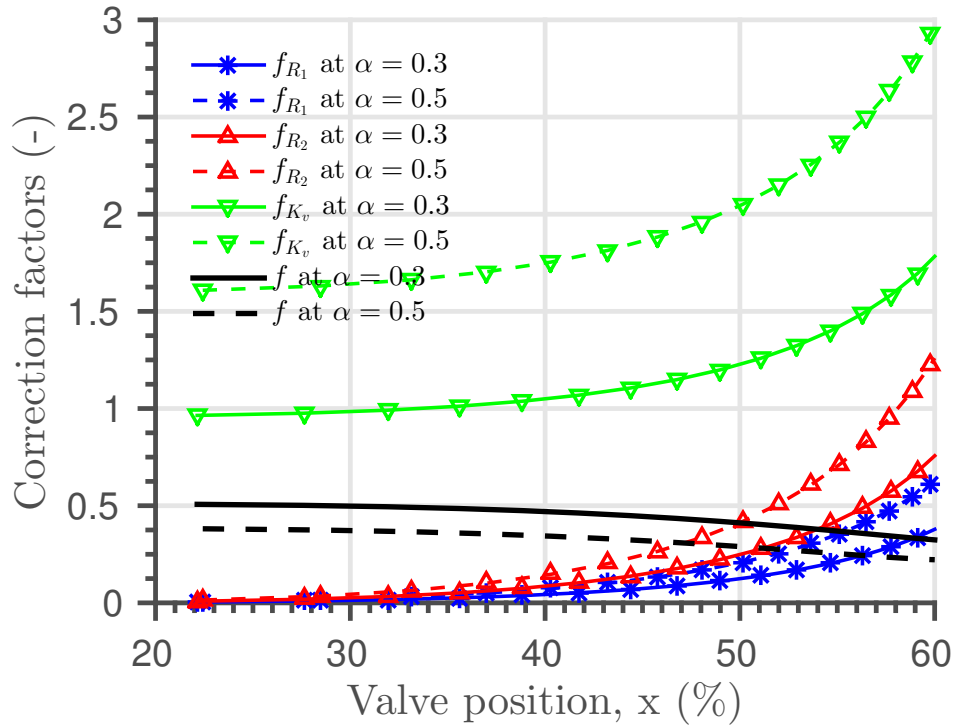


Fig. 12. Impact of R_1 , R_2 and K_v on the scaling function $f(\mathbf{p})$ between $K_{connected}(x)$ and $K_{isolated}(x)$.

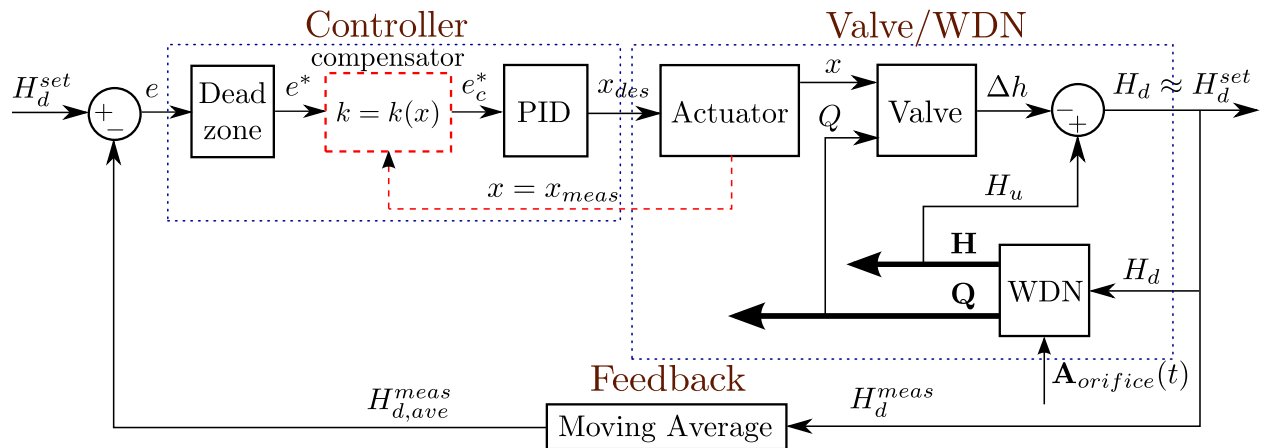


Fig. 13. Valve/WDN closed-loop model structure adopted from real-life system and implemented in the mathematical model used for the transient simulation study.

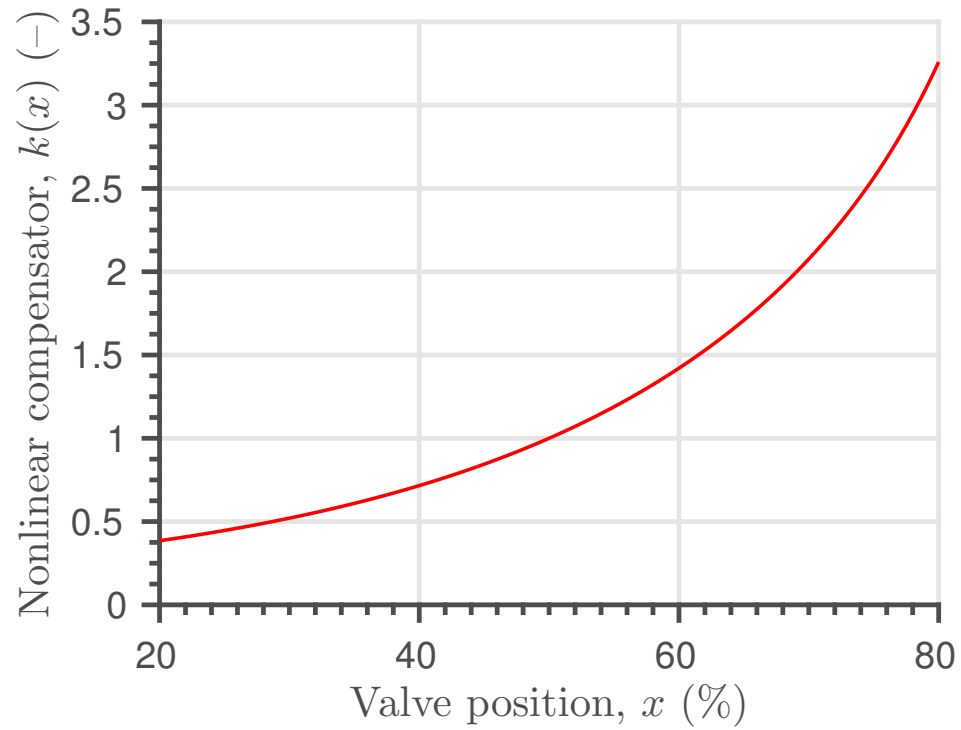


Fig. 14. Nonlinear gain compensator curve $k(x)$ used for scaling the Valve/WDN plant gain K in the transient simulation study.

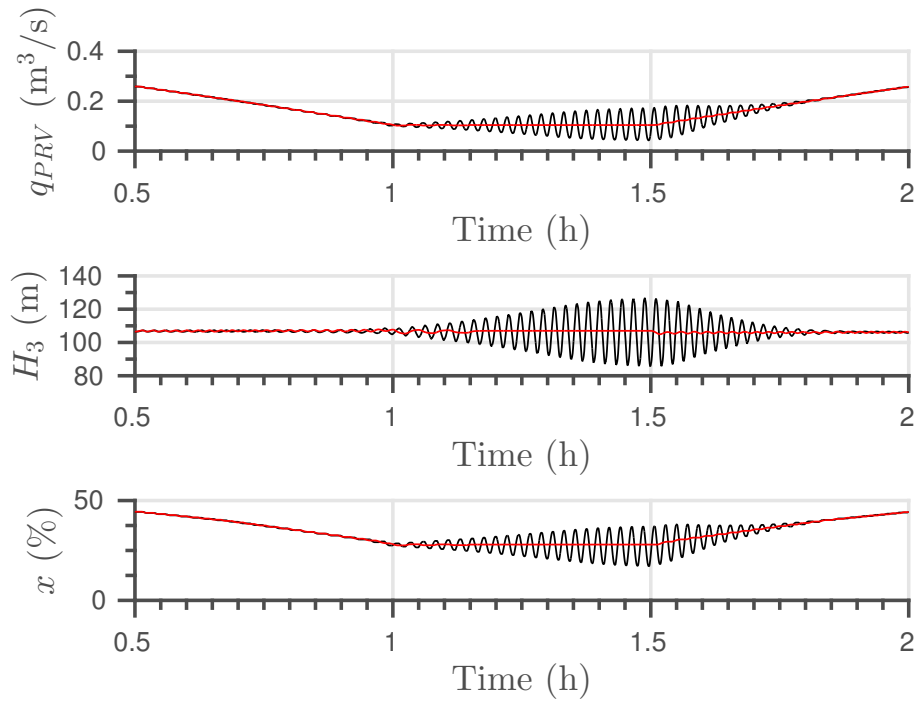


Fig. 15. Results of the transient simulations of our simplified pipe-PRV-pipe hydraulic model with (red) and without (black) gain compensation.